

What happens when the Damping Factor of an Audio-Amplifier is changed?

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This study researches the influence of the damping factor of an audio amplifier on the 'quality' of sound reproduction. Large changes are noticed and their causes and consequences are described and measured in detail.

Introduction

Repeatedly I observed that low damping factor valve amplifiers sound 'better' than high damping factor transistor amps. Why this is the case? It became clear that the observed differences are not an issue of 'valves versus transistors'. A transistor amp can sound 'better' when its damping factor is adjusted.

'It sounds better' is a very subjective observation. In this research I will bring observations from the subjective to the objective domain. Valves versus transistors will not be dealt with. The only focus is at the influence of the damping factor (DF) on the perceived sound quality and its explanation.

The research is organized as follows:

Chapter one describes research-amplifiers with adjustable DF.

Chapter two researches how the output voltage of an amplifier is influenced by its DF.

Chapter three discusses the transfer from amplifiers input voltage to perceived sound pressure levels.

Chapter four studies the amplifiers output current through the speakers voice coil.

Chapter five studies the acoustical distortions under different DF conditions.

Chapter six describes a controlled 'subjective' test where voltage drive (high DF) and current drive (low DF) are compared.

Chapters seven and eight give conclusions about 'sound quality' in relation to the DF of the applied audio amplifier.

1 : Research Amplifiers with adjustable DF

I designed several methods to change the DF of an amplifier. I applied the Hypex Ncore 122mp 2 x 120 W class-D amplifier.

1-a : Series Resistor

in series with the amplifiers output an extra resistor R-series is placed. The actual output impedance is enlarged by this series resistor: $Z_{out} = Z_{out-122mp} + R_{series}$.

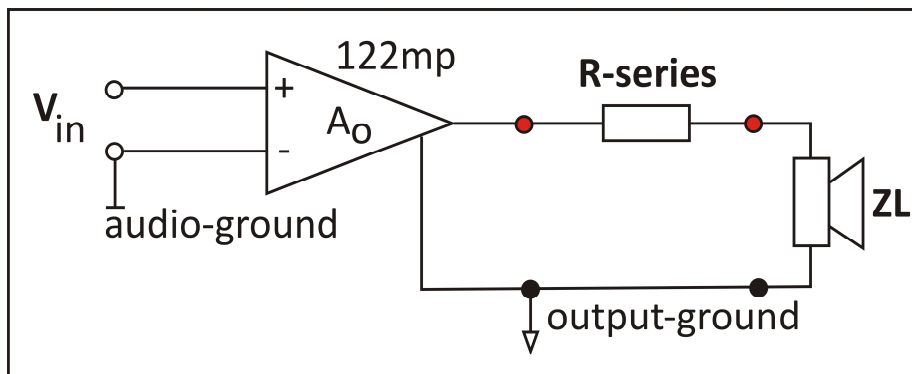


Figure 1: Z_{out} is enlarged by R-series

The DF is given by $DF = Z_L(f) / Z_{out}$ where $Z_L(f)$ is the impedance of the applied loudspeaker. Mostly one takes for $Z_L(f)$ a pure resistive load, say 4 or 8 Ohms. We can write: $DF_4 = 4 / Z_{out}$ or $DF_8 = 8 / Z_{out}$.

For the rest of this research I shall not use the DF, but refer directly to the effective Z_{out} . Then the choice of $Z_L(f)$ and its dependency of the frequency becomes arbitrary.

1-b : External Active Variable Z_{out} circuit

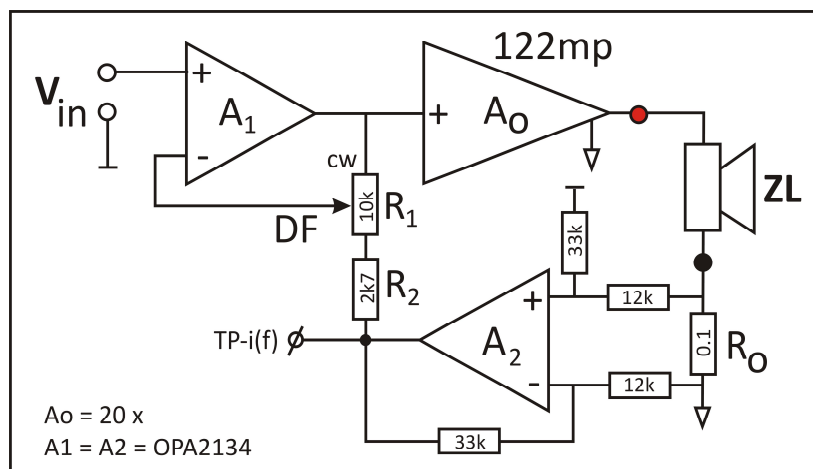


Figure 2: Active circuit around Nc-122mp creates adjustable Z_{out} .

In figure-2 the loudspeaker is connected to ground via a small resistor R_0 (use a 5 Watt type). The current through the loudspeaker generates a voltage over R_0 . This voltage is amplified (A_2) and added inverted to the input amp A_1 , which output is fed to the input of the Nc-122mp amp. Nothing is changed inside this class-D amp.

With the components shown, the actual output impedance of the complete circuit can be changed from 0.1Ω to 18Ω . Make R_2 smaller for a larger Z_{out} -value, but watch stability! I used a differential amp A_2 over R_0 to deal well with grounding-differences (audio-ground is unequal output-ground).

For correct functioning of this circuit in the frequency domain, drastic measures are taken to ensure stability, to correct for phase shift, because the Nc-122mp behaves as a 2-nd order low-pass. These measures are not drawn in the circuit. I have found a stable solution, still open for further optimization, but this remains outside the research.

Equation-3 shows the condition for constant V_{out} , independent of the wiper position of pot R_1 . I designed for $Z_L(f) = 5 \Omega$ at 320 Hz. Adapt A_2 for any other loudspeaker.

$$\text{Equation 3: } A_2 = (Z_L + R_0) / (A_0 \cdot R_0)$$

1-c : Add series resistor and change the FB-circuit inside an amplifier

When you have access to the feedback-circuit of an amplifier, without destroying the amp or killing its stability, you can use the next circuit.

I do not further discuss or use this solution that I designed in 1974. It works great, and the principle is: the weakening of amplification, caused by the current feedback, is compensated by enlarging the effective amplification of the amplifier, by changing its feedback ratio.

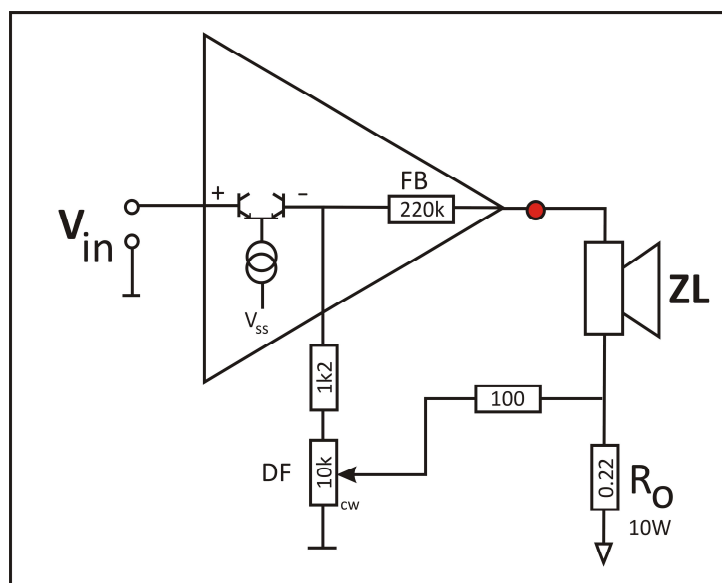


Figure 4: Z_{out} change combined with FB-change, resulting in constant V_{out} , independent of the DF-wiper-position.

2: $Z_L(f)$ creates a non-constant and more distorted $V_{out}(f)$ for $Z_{out} > 0$

I measured the output impedance and phase (circuit 1-b) for $Z_{out}(f) = 18 \Omega$. At highest frequencies some decline is noticeable, caused by the not yet optimal frequency corrections as mentioned in chapter 1-b. Also the electrical phase is shown. The measurements up to 20 kHz for $Z_{out} = 0,1$ and $Z_{out} = 2 \Omega$ just show straight lines.

This is the first indication that methods 1-a and 1-b behave almost equal, while method 1-b creates negligible power loss, and is therefore favored by me.

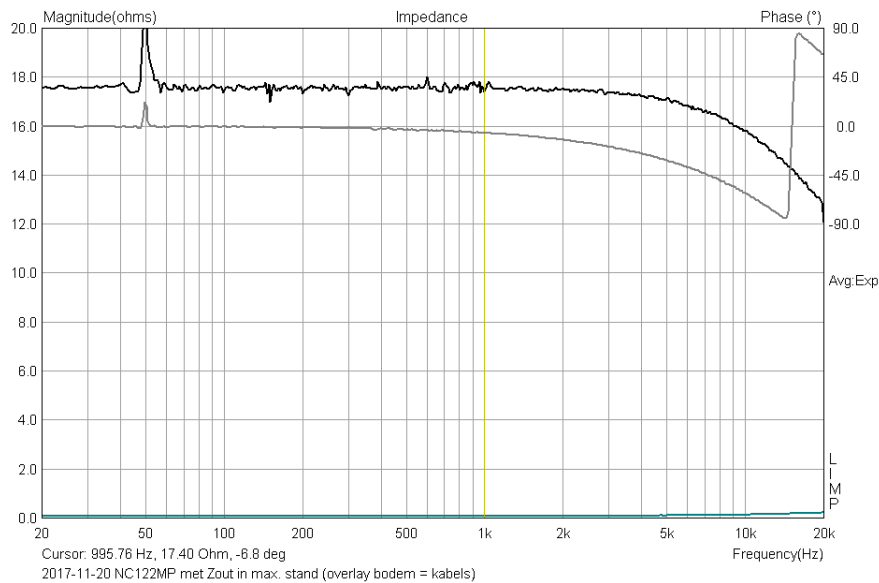


Figure 5: $Z_{out}(f)$ of circuit 1-b for the $Z_{out} = 18 \text{ Ohm}$ setting.

The electrical phase is shown by the light-grey curve (right hand scale).

The impedance of my test loudspeaker as function of frequency is shown below. The upper curve is its electrical phase.

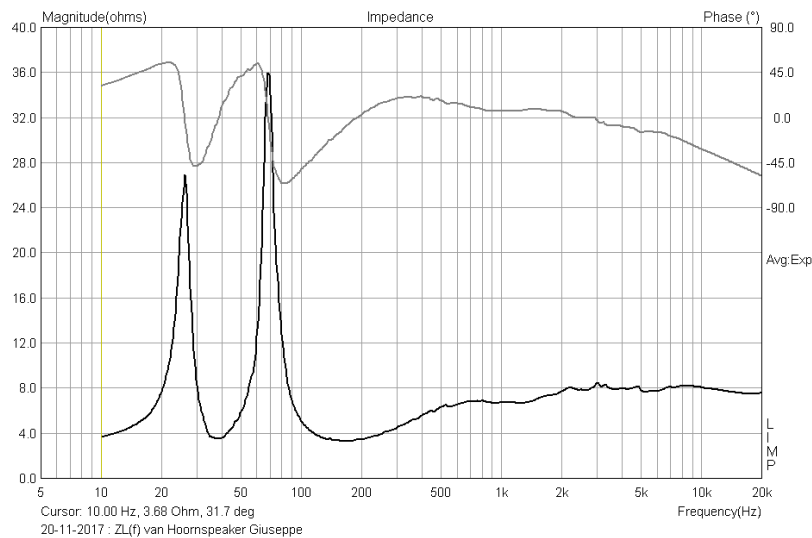


Figure 6: Electrical impedance and phase of the applied test loudspeaker

The next measurement shows V_{out} (over the loudspeaker) divided by V_{in} (amp input voltage) expressed in dBV/V = $[20\log(V_{out}/V_{in})]$ for three different values of Z_{out} : 0,1 and 2 and 18 Ω .

For $Z_{out} = 18 \Omega$ (red curve) above 20 kHz, some deviation is visible. See my earlier remark in chapter 1 about the not yet optimal frequency correction.

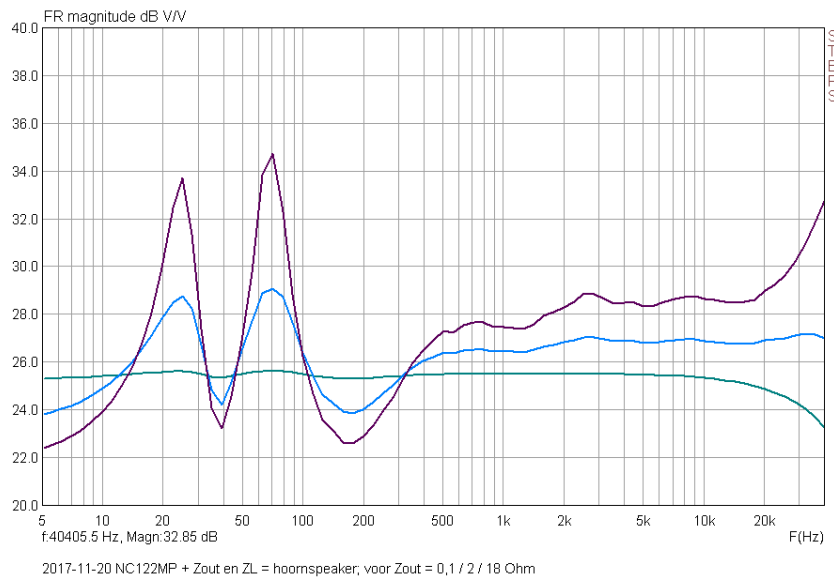


Figure 7: Effective amplification for $Z_{out} = 0,1$ (green) and 2 (blue) and 18 (red) Ω .

These results are fully described by $A = A_0 \cdot Z_L(f) / (Z_{out} + Z_L(f))$. I repeated this measurement with circuit 1-a and got exactly the same results. Again the equality of methods 1-a and 1-b is underlined.

Before delving into the huge consequences of this measurement, I also checked what happens with the voltage distortion. The current through the loudspeaker is a complex current, in which resonances and so called “back-emf” effects occur. The actual voltage over the speaker terminals is not only determined by the voltage amplification of the amplifier, but also by the changed current through the loudspeaker.

The next measurements show the distortion of the voltage over the speaker (at +6dBV level at 320 Hz) for $Z_{out} = 0,1$ and 18 Ω . Circuits 1-a and 1-b gave the same results.

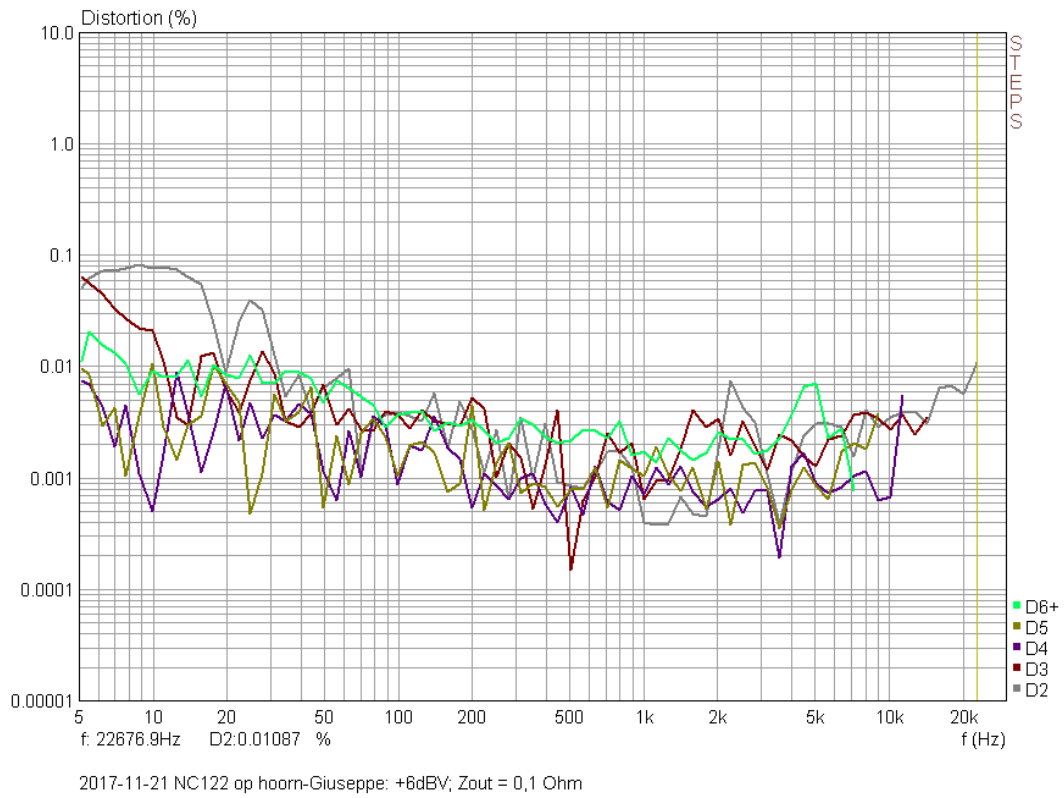


Figure 8: Vout distortion for Zout = 0,1 Ω

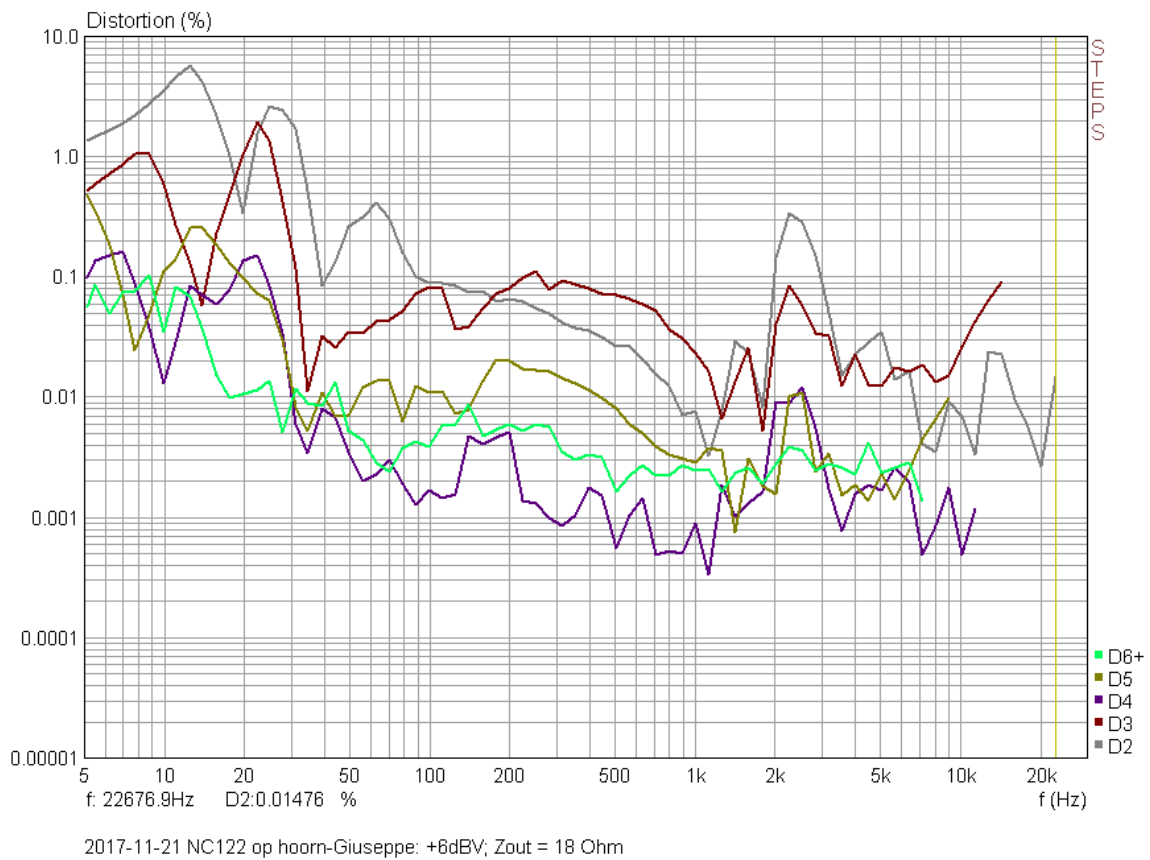


Figure 9: Vout distortion for Zout = 18 Ω

It is interesting to check the distortion resonance frequencies for Harmonic-2 and Harmonic-3 and to compare them with the resonance frequencies of the loudspeaker impedance (figure 6). Counting “1”, “2”, “3” helps a lot. Around 2,5 kHz a resonance in the mid-horn section is visible. With music I had not yet detected this by ear. Using a swept sine wave made it noticeable.

Let me draw some initial conclusions:

Conclusion-1:

Changing Z_{out} can equally be done with circuits 1-a and 1-b. These circuits behave equal, except for the larger power loss in circuit 1-a. Phase effects are not yet studied.

Conclusion-2:

Raising Z_{out} creates a $V_{out}(f)$ over the speaker terminals which follows the $Z_L(f)$ characteristics.

Conclusion-3:

Raising Z_{out} introduces distortions in $V_{out}(f)$ if $Z_L(f)$ is a standard loudspeaker which impedance is not corrected (see acknowledgement ‘Hans’ for more details).

With these results, one might say: “Transistor amps with infinitive damping are ideal, because they give constant V_{out} with no voltage distortion”. But why do such ideal amps often sound flat, no musicality, no enveloping,”? I need some more research to answer this repeated observation.

3: Measurements in the Acoustical domain

Up till now all measurements were in the electrical domain. Now I go acoustical by means of a mike and measure the resulting sound pressure levels.

In the electrical domain I already have measured the difference in transfers between the $Z_{out}=18$ and the $Z_{out}=0,1$ situations. See the figure below, compare to figure 7:

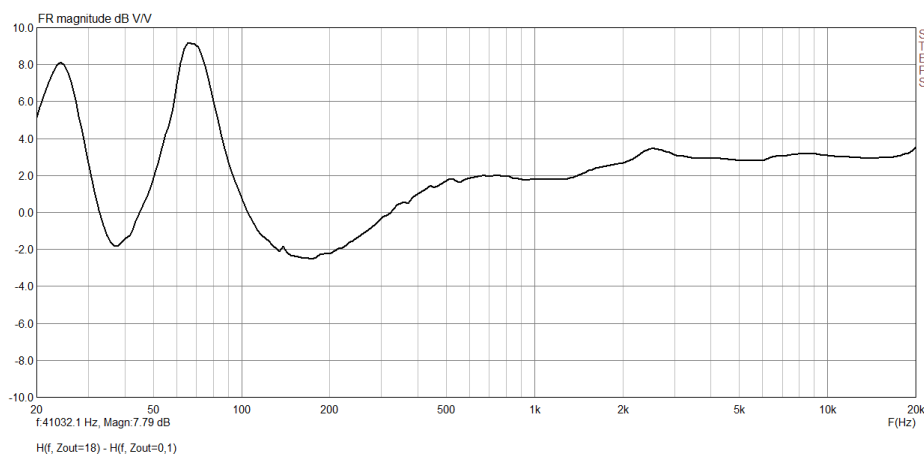


Figure 10: difference in electrical transfer between $Z_{out}=18$ and $Z_{out}=0,1$

In the acoustical domain I measured the transfers for $Z_{out}=18$ and $Z_{out}=0,1$ and determined their difference. See figure 11.

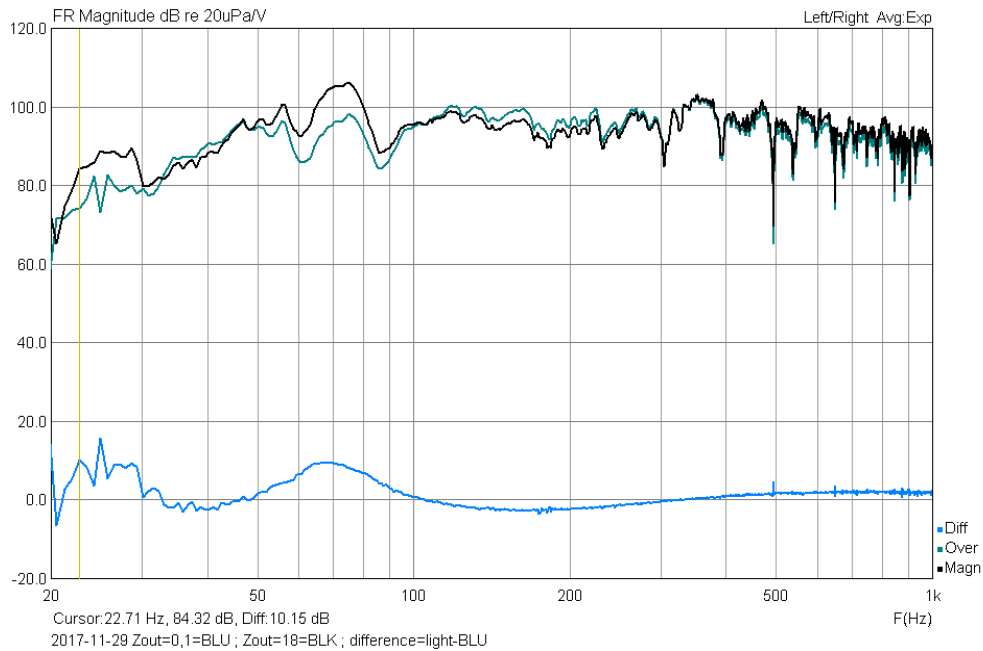


Figure 11: Acoustical transfers for $Z_{out}=18$ (black) and $Z_{out}=0,1$ (green) and their difference (blue bottom curve)

One clearly can see that the acoustical transfer-difference follows the electrical one (compare to figure 10).

Conclusion-4:

The electrical and acoustical transfers show the same behavior with respect to Z_{out} .

4 : $Z_L(f)$ creates a more constant and less distorted current $i(f)$ for $Z_{out} > 0$

'Rob' (see acknowledgement) pointed out that also the current through the loudspeaker should be studied. What actually drives the loudspeaker cone is the magnetic force $F=B \cdot i \cdot l$ which directly generates the acoustical sound pressure P_{spl} . Here F is the magnetic driving force to the cone, B = the effective flux density in the voice coil gap, i = the current through the voice coil and l = the effective length of the voice coil wire that sees B . So, at first sight, it is only the current through the voice coil that generates the 'music'.

See figure 12 for the currents through the voice coil, measured at the output of A2 (see figure 2) for $Z_{out} = 0,1$ and 18Ω .

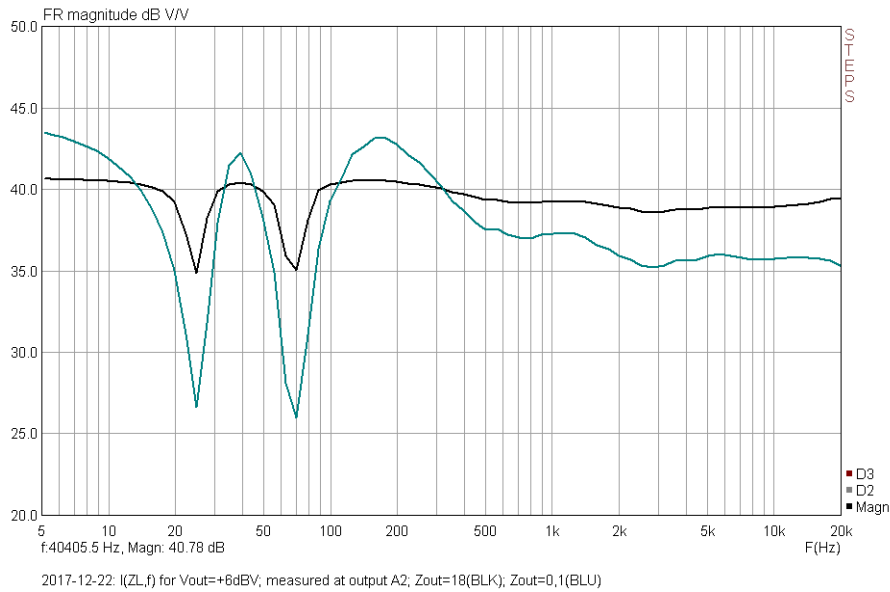


Figure 12: Current through the voice coil for Zout is 0,1 (blue) and 18 (black).
 Vertical scale conversion factor: 48 dBV/V = 1 A_{rms} through the voice coil at 320 Hz.

Indeed, if Zout raises in value, the current stays more constant. This is as expected and fully described by $A = V_{out}/V_{in} = A_0 \cdot Z_L(f) / (Z_{out} + Z_L(f))$ with $i(f) = V_{out} / Z_L(f)$.

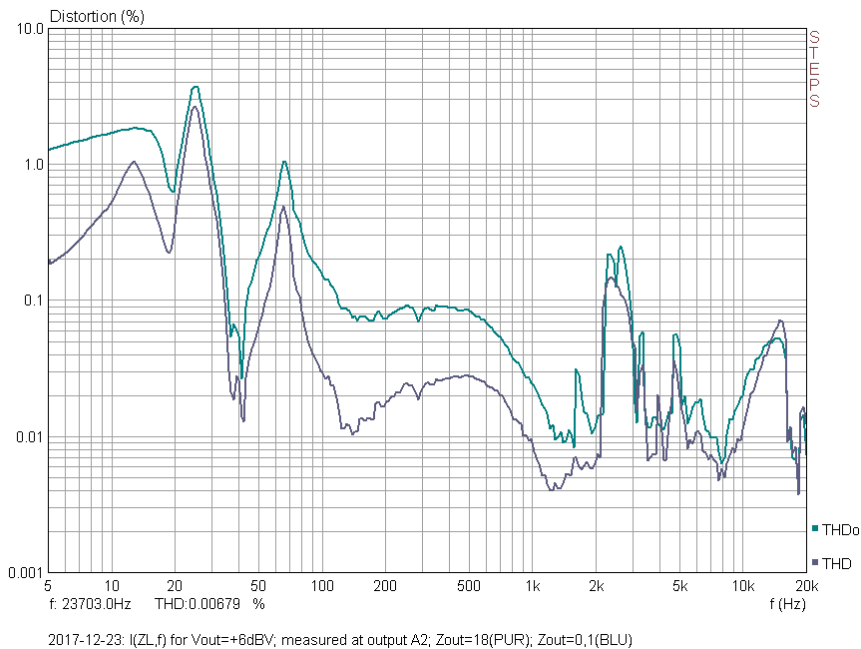


Figure 13: Total Harmonic Distortion of the current $i(f)$ through the voice coil for Zout=18 (grey) and Zout=0,1 (blue).

Figure 13 shows that the current gets less distorted when Zout raises. This means that the raise in voltage distortion (as measured in chapter 2) is a logical consequence to make $i(f)$ less distorted. Again this is a direct result of applying current feedback.

Conclusion-5 :

Enlarging Zout reduces the harmonic distortion in the current through the voice coil.

5 : Distortions in the Acoustical Domain

In chapter 3 the Z_{out} -effect on V_{out} was researched with the acoustical transfer. In this chapter I investigate the $i(f)$ -effects by measuring the acoustical distortion with an omni pressure microphone.

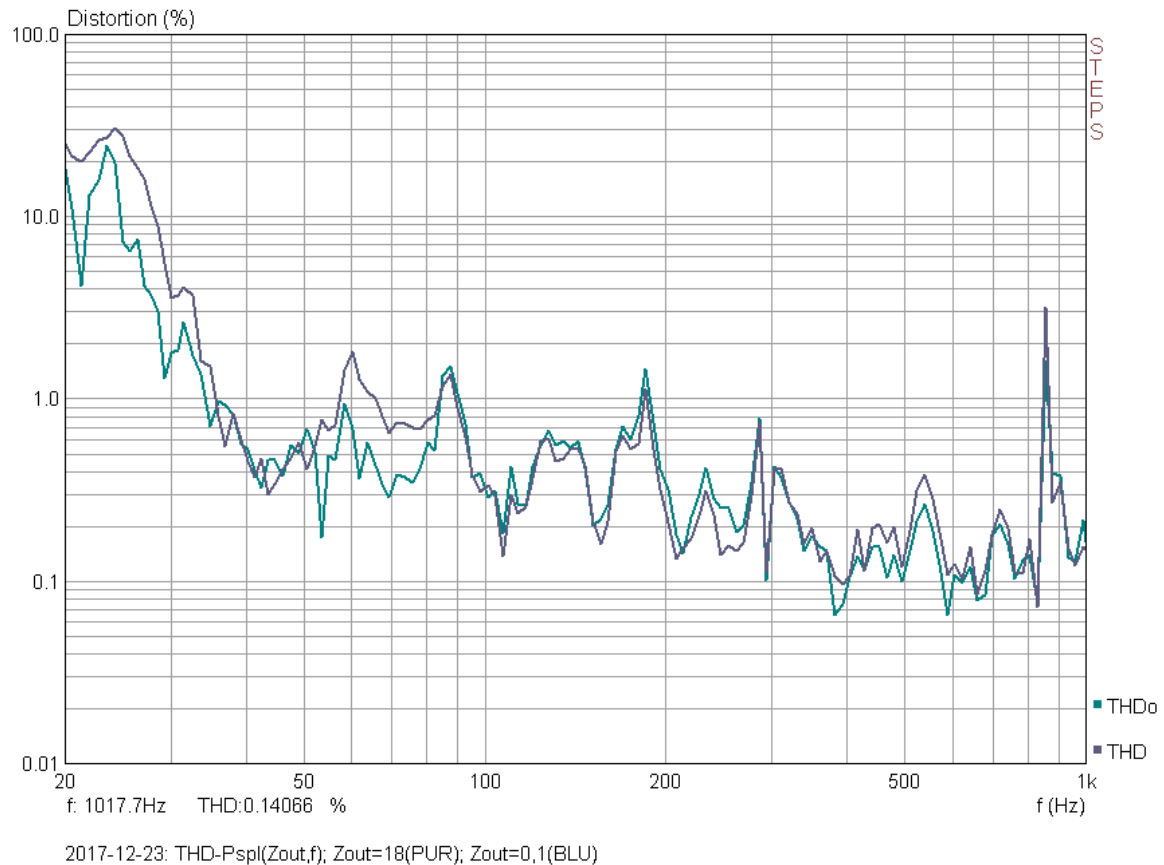


Figure 14: THD of the acoustical P_{spl} for $Z_{out}=0.1$ (blue) and 18 (grey).

This result is unexpected. If the current drives the voice coil, and if this current is less distorted for $Z_{out}=18$, then P_{spl} should be less distorted as well, which is not the case.

This measurement clearly shows that other distortions in the loudspeaker or environment surpass the electrical distortion (see figure 13). Reflections and room resonances are visible. With my gated-mike-swept-sine measurement I can't improve any further.

Therefore I compared my measurements with the results of 'Rob' (see acknowledgements). He applied an accelerometer mounted on the cone of a woofer for voltage drive (blue, $Z_{out}=0$) or current drive (red, $Z_{out}>100$).

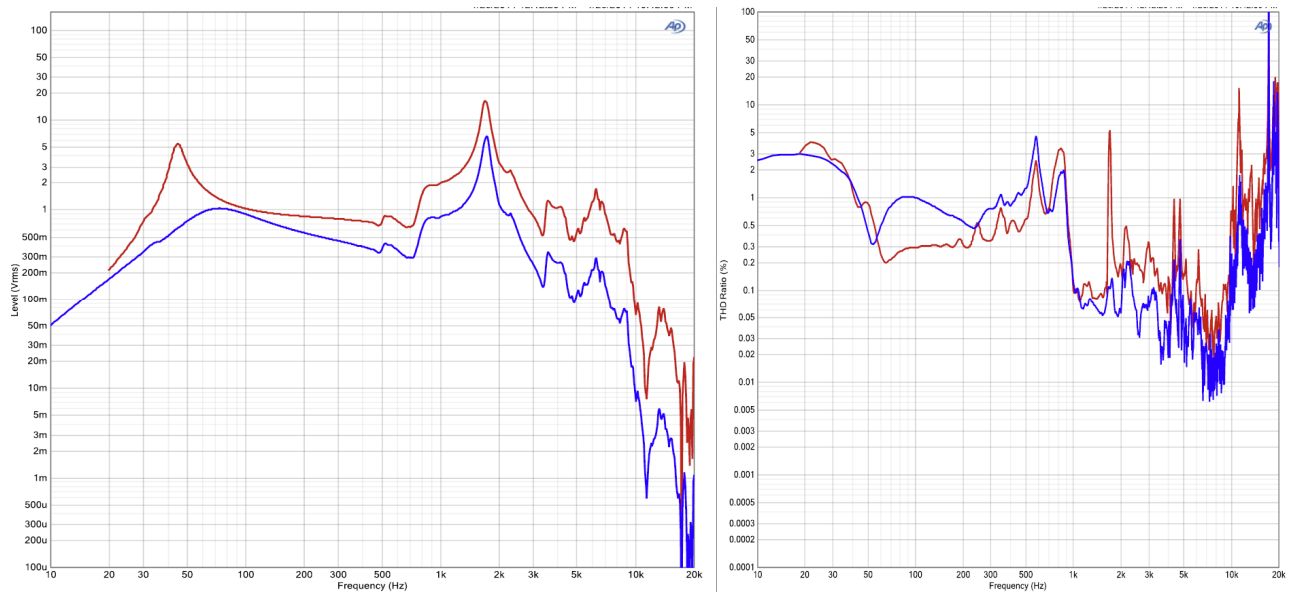


Figure 15: Source: <http://rmsacoustics.nl/activecontrol.html>

blue = voltage amplifier; red = current amplifier.

Output taken with an accelerometer mounted on a woofer cone.

Reflections and disturbances have less influence in measurement 15. From 50 Hz to 500 Hz the acoustical distortion (right figure) is less with large Z_{out} amps.

One might argue that I can't hear under the distortion level as measured in figure 14. I doubt such line of reasoning. First, my ears do not function as a gated-swept-sine-measurement. Secondly, with my loudspeakers, I can distinguish clearly if an amplifier distorts 0.3 %THD or 0.03 %THD. I even am able to notice if in 0.03 %THD the second or third harmonic is dominant. My ear is far more better than the measurement method applied in figure 14. So, I am fully confident to use the measurement results of 'Rob' and to consider them as reliable.

Conclusion-6 :

Current drive (Z_{out} is large) gives less acoustical distortion compared to voltage drive (Z_{out} is small). The frequency region where this is the case is speaker dependent. In my research the sound between 100 Hz and 2 kHz gets less distorted.

6: Lets go Subjective

Two subjective test are described. First: how does the sound character change for $Z_{out} = 18$ compared to $Z_{out} = 0,1 \Omega$. Secondly: if we change a reproduced recording with the electrical voltage ($Z_{out}=18 / 0,1$) difference curve, are we then able to notice any difference?

6-1: Differences in Sound Character.

I took a piece of music of Eva Cassidy and listened to it under the $Z_{out}=18$ and $Z_{out}=0,1$ conditions at equal loudness (at 320 Hz). Compared to the $Z_{out}=0,1$ situation the sound character of the $Z_{out}=18$ was:

- a) More brightness
- b) More deep base
- c) More mechanical vibrations of my lab-table
- d) More sense of being enveloped
- e) More details noticeable

Let's be honest: observations a) and b) are totally determined by the changed voltage transfer function. Electrically there is more voltage at higher frequencies and also around the low frequency speaker resonances. These observations are totally understandable.

Observation c) deals with the lowest resonance frequency (around 25 Hz) which I do not hear, but sense as table vibrations. With $Z_{out}=18$ there is more V_{out} and therefore this region is reproduced louder, creating more table vibrations.

Observation d) en e): Might it be true that I happen to like more base and brightness and that I express that by 'envelopment' and 'details'? I will deal with this assumption and more in the next sections.

6-2: Recording corrected with the electrical voltage-transfer-difference curve

The next test was a little more difficult to do. I took the Eva Cassidy recording and multiplied it with the electrical voltage transfer difference between $Z_{out}=18$ and $Z_{out}=0,1$ (figure 10).

I listened to this changed recording with the amplifier in the $Z_{out}=0,1$ setting (changed-Eva- $Z_{out}=0,1$).

Following I listened to the unchanged recording with the amp in $Z_{out}=18$ setting (unchanged-Eva- $Z_{out}=18$) and compared many times with the (changed-Eva- $Z_{out}=0,1$) reproduction.

I made the following observations:

- a) Equal brightness
- b) Equal base
- c) The low-mid part is more 'natural/honest/open' in unchanged-Eva- $Z_{out}=18$

Observations a) and b) are explainable in the electrical domain. Observation c) is more hidden. In the next chapter I will discuss it.

7 : Discussions

As said: observations a) and b) can be explained with the now compensated voltage changes. In my opinion there is no need for further discussion. That is the part we all fully understand. I might like more high and base, but in this last test these levels are equal. So my possible preference is pushed aside.

Let us now focus on observation c). By multiplying “Eva” with the electrical difference transfer, the result is a little larger V_{out} . Because the Nc-mp122 amplifier is extremely clean, negligible extra voltage distortion is introduced. Therefore the reproduction with $Z_{out}=0,1$ does not contain extra V_{out} distortion. With $Z_{out}=18$, the previous chapter has shown that the acoustical distortion is less between 100 Hz and 2 kHz. Meaning: the two test situations of chapter 6 have different acoustical distortion artifacts. Situation unchanged-Eva- $Z_{out}=18$ is the cleanest in the “low-mid”-region, which I clearly observed.

8 : Conclusions

This research has shown that raising the output impedance of an amplifier, generates changes in the electrical transfer, based on the frequency dependence of the impedance of the applied loudspeaker. The changed transfer function in the electrical domain almost equally appears in the acoustical domain. With infinitive damping ($Z_{out} = 0 \text{ Ohm}$) the amplifiers output voltage is only distorted by the amplifier itself. With $Z_{out} > 0$, the current feedback generates extra distortions in the voltage over the loudspeaker. However, when we consider the current through the loudspeaker, this current gets less distorted for $Z_{out} > 0$. Because, in first approximation, it is the current that moves the loudspeaker cone. The acoustical distortion of the speaker gets less for $Z_{out} > 0$. In such a situation the acoustical reproduction is cleaner and this explains why it is preferred by the author.

9 : Acknowledgement

I got much support from the persons mentioned following. I thank all deeply for taking the time to study my attempts to find answers and to give excellent advises.

Rob Munnig Schmidt (RMSacoustics.nl); Guido Tent (Tentlabs.com and GrimmAudio.com); Hans van Maanen (TemporalCoherence.nl); Jan Didden (LinearAudio.net); Ari Polisois (polisois-audio.com); Pierre Touzelet (see AudioXpress.com and LinearAudio.net); Joe Rasmussen (custumanalogue.com).

10 : Literature

See the websites indicated in the Acknowledgement to find many high quality references.

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The Netherlands; Hichtum; Menno van der Veen ; www.mennovanderveen.nl